In this analysis, scattering amplitude can be written as

$$f(\theta) = \frac{1}{2ik} \sum_{\ell=0}^{\infty} (2\ell+1) \left[e^{2i\delta_{\ell}} - 1 \right] P_{\ell}(\cos\theta)$$
$$\Rightarrow f(\theta) = \frac{1}{k} \sum_{\ell=0}^{\infty} (2\ell+1) e^{i\delta_{\ell}} \sin\delta_{\ell} P_{\ell}(\cos\theta)$$

where $k^2 = \frac{2\mu E}{\hbar^2}$ and δ_{ℓ} is the phase shift of the individual partial waves due to scattering.

The corresponding differential cross-section will be

$$\frac{d\sigma}{d\theta} = \left| f\left(\theta\right)^2 \right| = \frac{1}{k^2} \left| \sum_{\ell=0}^{\infty} (2\ell+1) e^{i\delta_{\ell}} \sin \delta_{\ell} P_{\ell} \left(\cos \theta\right) \right|^2$$

Therefore, the total cross-section will be

$$\sigma_t = \int_0^{\pi} \int_0^{2\pi} \frac{d\sigma}{d\Omega} \sin\theta \, d\theta \, d\phi = 2\pi \int_0^{\pi} \left| f\left(\theta, \phi\right) \right|^2 \sin\theta \, d\theta \qquad = \frac{4\pi}{k^2} \sum_{\ell=0}^{\infty} (2\ell+1) \sin^2 \delta_\ell$$

Example: For *S* waves $(\ell = 0)$,

Scattering amplitude
$$f(\theta) = \frac{1}{k} e^{i\delta_0} \sin \delta_0$$

Differential cross section $\frac{d\sigma}{d\Omega} = \left| f(\theta) \right|^2 = \frac{1}{k^2} \sin^2 \delta_0$

Total cross section
$$\sigma = \frac{4\pi}{k^2} \sin^2 \delta_0$$

Note:

(i) Total cross section does not depend upon on the scattering angle θ .

(ii) At sufficiently low energies, s-wave scattering will be dominating one. As energy increases, contributions from the higher value of *l* values become important.

Optical Theorem

Now, for $\theta = 0$

$$f(0) = \frac{1}{k} \sum_{\ell=0}^{\infty} (2\ell+1) e^{i\delta_{\ell}} \sin \delta_{\ell} = \frac{1}{k} \sum_{\ell=0}^{\infty} (2\ell+1) \left[\cos \delta_{\ell} \sin \delta_{\ell} + i \sin^{2} \delta_{\ell} \right]$$
$$\sigma_{t} = \frac{4\pi}{k^{2}} \sum_{\ell=0}^{\infty} (2\ell+1) \sin^{2} \delta_{\ell} = \frac{4\pi}{k} \operatorname{Im} \left[f(0) \right]$$

Now,

This relation is known as Optical theorem. **Phase shift:**

$$\sin \delta_{\ell} = -\frac{2\mu k}{\hbar^2} \int_0^\infty V(r) j_l^2(kr) r^2 dr$$

For attractive potential $V(r) = -ve \implies$ positive phase shift. For repulsive potential $V(r) = +ve \implies$ negative phase shift.

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