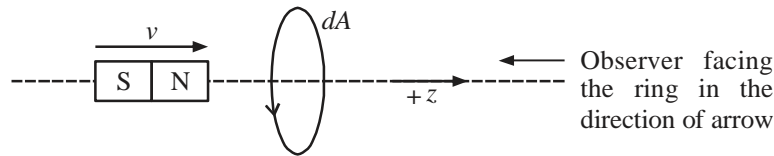
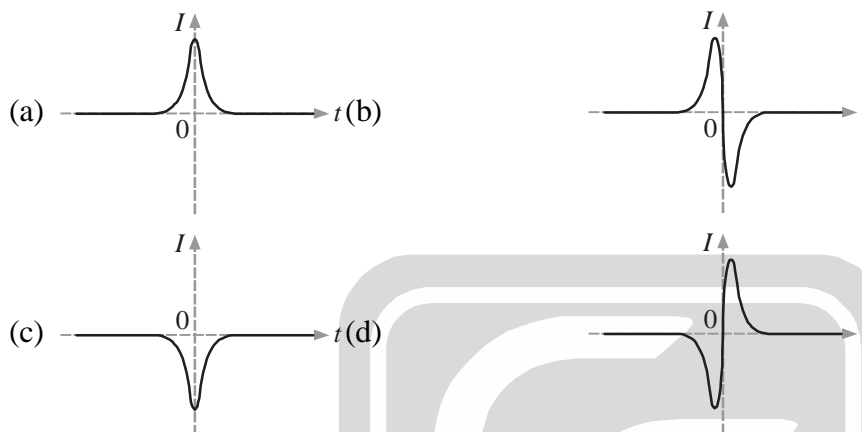


12. The ratio C_P/C_V of the specific heats at constant pressure and volume of a monoatomic ideal gas in two dimension is
 (a) $3/2$ (b) 2 (c) $5/3$ (d) $5/2$
13. A discrete random variable X takes a value from the set $\{-1, 0, 1, 2\}$ with the corresponding probabilities $p(X) = \frac{3}{10}, \frac{2}{10}, \frac{2}{10}$ and $\frac{3}{10}$, respectively. The probability distribution $q(Y) = (q(0), q(1), q(4))$ of the random variable $Y = X^2$ is
 (a) $\left(\frac{1}{5}, \frac{3}{5}, \frac{1}{5}\right)$ (b) $\left(\frac{1}{5}, \frac{1}{2}, \frac{3}{10}\right)$ (c) $\left(\frac{2}{5}, \frac{2}{5}, \frac{1}{5}\right)$ (d) $\left(\frac{3}{10}, \frac{3}{10}, \frac{2}{5}\right)$
14. The total number of phonon modes in a solid of volume V is $\int_0^{\omega_D} g(\omega) d\omega = 3N$, where N is the number of primitive cells, ω_D is the Debye frequency and density of photon modes is $g(\omega) = AV\omega^2$ (with $A > 0$ a constant). If the density of the solid doubles in a phase transition, the Debye temperature θ_D will
 (a) increase by a factor of $2^{2/3}$ (b) increase by a factor of $2^{1/3}$
 (c) decrease by a factor of $2^{2/3}$ (d) decrease by a factor of $2^{1/3}$
15. The volume of the region common to the interiors of two infinitely long cylinders defined by $x^2 + y^2 = 25$ and $x^2 + 4z^2 = 25$ is best approximated by
 (a) 225 (b) 333 (c) 423 (d) 625
16. In an experiment to measure the charge to mass ratio e/m of the electron by Thomson's method, the value of the deflecting electric field and the accelerating potential are 6×10^6 N/C (Newton per coulomb) and 150 V, respectively. The magnitude of the magnetic field that leads to zero deflection of the electron beam is closest to
 (a) 0.6 T (b) 1.2 T (c) 0.4 T (d) 0.8 T
17. The volume and temperature of a spherical cavity filled with black body radiation are V and 300 K, respectively. If it expands adiabatically to a volume $2V$, its temperature will be closest to
 (a) 150 K (b) 300 K (c) 250 K (d) 240 K
18. A particle, thrown with a speed v from the earth's surface, attains a maximum height h (measured from the surface of the earth). If v is half the escape velocity and R denotes the radius of earth, then h/R is
 (a) $2/3$ (b) $1/3$ (c) $1/4$ (d) $1/2$
19. The vector potential for an almost point like magnetic dipole located at the origin is $A = \frac{\mu \sin \theta}{4\pi r^2} \hat{\phi}$, where (r, θ, ϕ) denote the spherical polar coordinates and $\hat{\phi}$ is the unit vector along ϕ . A particle of mass m and charge q , moving in the equatorial plane of the dipole, starts at time $t = 0$ with an initial speed v_0 and an impact parameter b . Its instantaneous speed at the point of closest approach is
 (a) v_0 (b) 0/0 (c) $v_0 + \frac{\mu q}{4\pi m b^2}$ (d) $\sqrt{v_0^2 + \left(\frac{\mu q}{4\pi m b^2}\right)^2}$

20. A conducting wire in the shape of a circle lies on the (x, y) -plane, with its centre at the origin. A bar magnet moves with a constant velocity towards the wire along the z -axis (as shown in the figure below).



We take $t = 0$ to be the instant at which the midpoint of the magnet is at the centre of the wire loop and the induced current to be positive when it is counter-clockwise as viewed by the observer facing the loop and the incoming magnet. In these conventions, the best schematic representation of the induced current $I(t)$ as a function of t , is



21. The components of the electric field, in a region of space devoid of any charge or current sources, are given to be $E_i = a_i + \sum_{j=1,2,3} b_{ij} x_j$, where a_i and b_{ij} are constants independent of the coordinates.

The number of independent components of the matrix b_{ij} is

- (a) 5 (b) 6 (c) 3 (d) 4
22. The equation of motion of a one-dimensional forced harmonic oscillator in the presence of a dissipative force is described by $\frac{d^2x}{dt^2} + 10\frac{dx}{dt} + 16x = 6te^{-8t} + 4t^2e^{-2t}$. The general form of the particular solution, in terms of constants A, B etc., is

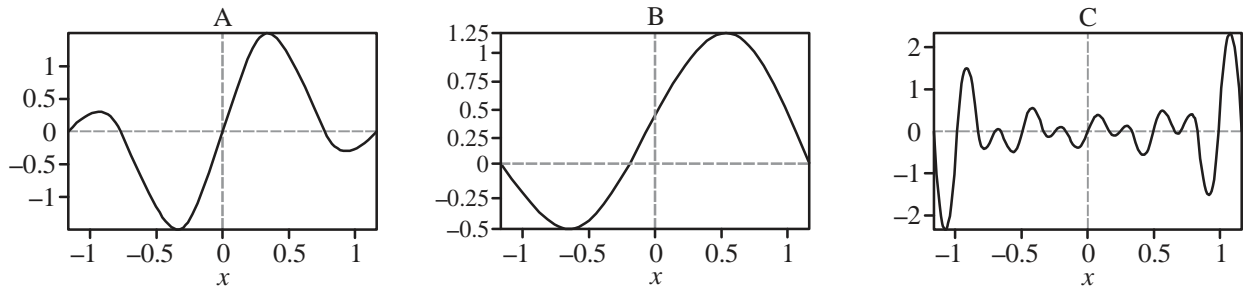
- (a) $t(At^2 + Bt + C)e^{-2t} + (Dt + E)e^{-8t}$ (b) $(At^2 + Bt + C)e^{-2t} + (Dt + E)e^{-8t}$
 (c) $t(At^2 + Bt + C)e^{-2t} + t(Dt + E)e^{-8t}$ (d) $(At^2 + Bt + C)e^{-2t} + t(Dt + E)e^{-8t}$

23. The position of a particle in one dimension changes in discrete steps. With each step it moves to the right, however, the length of the step is drawn from a uniform distribution from the interval $\left[\lambda - \frac{1}{2}w, \lambda + \frac{1}{2}w\right]$, where λ and w are positive constants. If X denotes the distance from the start-

ing point after N steps, the standard deviation $\sqrt{\langle X^2 \rangle - \langle X \rangle^2}$ for large values of N is

- (a) $\frac{\lambda}{2} \times \sqrt{N}$ (b) $\frac{\lambda}{2} \times \sqrt{\frac{N}{3}}$ (c) $\frac{w}{2} \times \sqrt{N}$ (d) $\frac{w}{2} \times \sqrt{\frac{N}{3}}$

24. The figures below depict three different wavefunctions of a particle confined to a one-dimensional box $-1 \leq x \leq 1$.



The wavefunctions that correspond to the maximum expectation values $|\langle x \rangle|$ (absolute value of the mean position) and $\langle x^2 \rangle$, respectively, are

- (a) B and C (b) B and A (c) C and B (d) A and B
25. The volume integral $I = \iiint_V \mathbf{A} \cdot (\nabla \times \mathbf{A}) d^3x$ is over a region V bounded by a surface Σ (an infinitesimal area element being $\hat{\mathbf{n}}dS$, where $\hat{\mathbf{n}}$ is the outward unit normal). If it changes to $I + \Delta I$, when the vector \mathbf{A} is changed to $\mathbf{A} + \nabla\Lambda$, then ΔI can be expressed as:

- (a) $\iiint_V \nabla \cdot (\nabla\Lambda \times \mathbf{A}) d^3x$ (b) $-\iiint_V \nabla^2\Lambda d^3x$
 (c) $-\oiint_{\Sigma} (\nabla\Lambda \times \mathbf{A}) \cdot \hat{\mathbf{n}}dS$ (d) $\oiint_{\Sigma} \nabla\Lambda \cdot \hat{\mathbf{n}}dS$

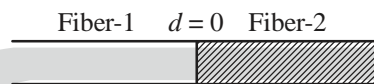
PART - C

1. A particle of mass m moves in a potential that is $V = \frac{1}{2}m(\omega_1^2 x^2 + \omega_2^2 y^2 + \omega_3^2 z^2)$ in the coordinates of a non-inertial frame F . The frame F is rotating with respect to an inertial frame with an angular velocity $\hat{\mathbf{k}}\Omega$, where $\hat{\mathbf{k}}$ is the unit vector along their common z -axis. The motion of the particle is unstable for all angular frequencies satisfying
- (a) $(\Omega^2 - \omega_1^2)(\Omega^2 - \omega_2^2) > 0$ (b) $(\Omega^2 - \omega_1^2)(\Omega^2 - \omega_2^2) < 0$
 (c) $(\Omega^2 - (\omega_1 + \omega_2)^2)(\Omega^2 - |\omega_1 - \omega_2|^2) > 0$ (d) $(\Omega^2 - (\omega_1 + \omega_2)^2)(\Omega^2 - |\omega_1 - \omega_2|^2) < 0$
2. Potassium chloride forms an FCC lattice, in which K and Cl occupy alternating sites. The density of KCl is 1.98 g/cm^3 and the atomic weights of K and Cl are 39.1 and 35.5, respectively. The angles of incidence (in degrees) for which Bragg peaks will appear when X-ray of wavelength 0.4 nm is shone on a KCl crystal are
- (a) 18.5, 39.4 and 72.2 (b) 19.5 and 41.9
 (c) 12.5, 25.7, 40.5 and 60.0 (d) 13.5, 27.8, 44.5 and 69.0
3. The energies of a two-state quantum system are E_0 and $E_0 + \alpha\hbar$ (where $\alpha > 0$ is a constant) and the corresponding normalized state vectors are $|0\rangle$ and $|1\rangle$, respectively. At time $t = 0$, when the system is in the state $|0\rangle$, the potential is altered by a time independent term V such that $\langle 1|V|0\rangle = (\hbar\alpha/10)$. The transition probability to the state $|1\rangle$ at times $t \ll (1/\alpha)$, is



(a) $\frac{\alpha^2 t^2}{25}$ (b) $\frac{\alpha^2 t^2}{50}$ (c) $\frac{\alpha^2 t^2}{100}$ (d) $\frac{\alpha^2 t^2}{200}$

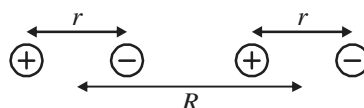
4. The Newton-Raphson method is to be used to determine the reciprocal of the number $x = 4$. If we start with the initial guess 0.20 then after the first iteration the reciprocal is
 (a) 0.23 (b) 0.24 (c) 0.25 (d) 0.26
5. A satellite of mass m orbits around earth in an elliptic trajectory of semi-major axis a . At a radial distance $r = r_0$, measured from the centre of the earth, the kinetic energy is equal to half the magnitude of the total energy. If M denotes the mass of the earth and the total energy is $-\frac{GMm}{2a}$, the value of r_0/a is nearest to
 (a) 1.33 (b) 1.48 (c) 1.25 (d) 1.67
6. A laser beam propagates from fiber 1 to fiber 2 in a cavity made up of two optical fibers (as shown in the figure). The loss factor of fiber 2 is 10 dB/km.



If $E_2(d)$ denotes the magnitude of the electric field in fiber 2 at a distance d from the interface, the ratio $\frac{E_2(0)}{E_2(d)}$ for $d = 10$ km, is

(a) 10^2 (b) 10^3 (c) 10^5 (d) 10^7

7. The $|3, 0, 0\rangle$ state (in the standard notation $|n, l, m\rangle$) of the H-atom in the non-relativistic theory decays to the state $|1, 0, 0\rangle$ via two dipole transitions. The transition route and the corresponding probability are
 (a) $|3, 0, 0\rangle \rightarrow |2, 1, -1\rangle \rightarrow |1, 0, 0\rangle$ and $\frac{1}{4}$ (b) $|3, 0, 0\rangle \rightarrow |2, 1, 1\rangle \rightarrow |1, 0, 0\rangle$ and $\frac{1}{4}$
 (c) $|3, 0, 0\rangle \rightarrow |2, 1, 0\rangle \rightarrow |1, 0, 0\rangle$ and $\frac{1}{3}$ (d) $|3, 0, 0\rangle \rightarrow |2, 1, 0\rangle \rightarrow |1, 0, 0\rangle$ and $\frac{2}{3}$
8. The dispersion relation of a gas of non-interacting bosons in d dimensions is $E(k) = ak^s$, where a and s are positive constants. Bose-Einstein condensation will occur for all values of
 (a) $d > s$ (b) $d + 2 > s > d - 2$
 (c) $s > 2$ independent of d (d) $d > 2$ independent of s
9. A linear diatomic molecule consists of two identical small electric dipoles with an equilibrium separation R , which is assumed to be a constant. Each dipole has charges $+q$ of mass m separated by r when the molecule is at equilibrium. Each dipole can execute simple harmonic motion of angular frequency ω .

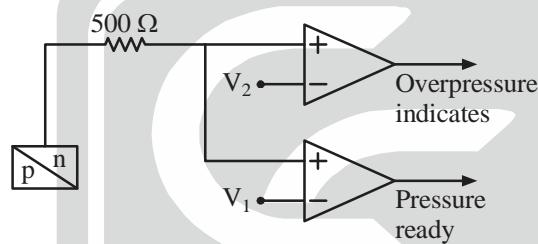


Recall that the interaction potential between two dipoles of moments \mathbf{p}_1 and \mathbf{p}_2 , separated by $\mathbf{R}_{12} = R_{12}\hat{n}$ is $(\mathbf{p}_1 \cdot \mathbf{p}_2 - 3(\mathbf{p}_1 \cdot \hat{n})(\mathbf{p}_2 \cdot \hat{n})) / (4\pi \epsilon_0 R_{12}^3)$.

Assume that $R \gg r$ and let $\Omega^2 = \frac{q^2}{4\pi \epsilon_0 mR^3}$. The angular frequencies of small oscillations of the diatomic molecule are

- (a) $\sqrt{\omega^2 + \Omega^2}$ and $\sqrt{\omega^2 - \Omega^2}$ (b) $\sqrt{\omega^2 + 3\Omega^2}$ and $\sqrt{\omega^2 - 3\Omega^2}$
 (c) $\sqrt{\omega^2 + 4\Omega^2}$ and $\sqrt{\omega^2 - 4\Omega^2}$ (d) $\sqrt{\omega^2 + 2\Omega^2}$ and $\sqrt{\omega^2 - 2\Omega^2}$

10. In the reaction $p + n \rightarrow p + K^+ + X$, mediated by strong interaction, the baryon number B , strangeness S and the third component of isospin I_3 of the particle X are, respectively.
 (a) $-1, -1$ and -1 (b) $+1, -1$ and -1 (c) $+1, -2$ and $-1/2$ (d) $-1, -1$ and 0
11. The pressure of a gas in a vessel needs be maintained between 1.5 bar to 2.5 bar in an experiment. The vessel is fitted with a pressure transducer that generates 4 mA to 20 mA current for pressure in the range 1 bar to 5 bar. The current output of the transducer has a linear dependence on the pressure.



The reference voltage V_1 and V_2 in the comparators in the circuit (shown in figure above) suitable for the desired operating conditions are, respectively

- (a) 2 V and 10 V (b) 2 V and 5 V (c) 3 V and 10 V (d) 3 V and 5 V
12. A ^{60}Co nucleus β -decays from its ground state with $J^p = 5^+$ to a state of ^{60}Ni with $J^p = 4^+$. From the angular momentum selection rules, the allowed values of the orbital angular momentum L and the total spin S of the electron-antineutrino pair are
 (a) $L = 0$ and $S = 1$ (b) $L = 1$ and $S = 0$ (c) $L = 0$ and $S = 0$ (d) $L = 1$ and $S = 1$
13. Diffuse hydrogen gas within a galaxy may be assumed to follow a Maxwell distribution at temperature 10^6 K, while the temperature appropriate for the H gas in the inter-galactic space, following the same distribution, may be taken to be 10^4 K. The ratio of thermal broadening $\Delta v_G / \Delta v_{IG}$ of the Lyman- α line from the H-atoms within the galaxy to that from the inter-galactic space is closest to
 (a) 100 (b) 1/100 (c) 10 (d) 1/10
14. A polymer made up of N monomers, is in thermal equilibrium at temperature T . Each monomer could be of length a or $2a$. The first contributes zero energy, while the second one contributes ϵ . The average length (in units of Na) of the polymer at temperature $T = (\epsilon/k_B)$ is
 (a) $\frac{5+e}{4+e}$ (b) $\frac{4+e}{3+e}$ (c) $\frac{3+e}{2+e}$ (d) $\frac{2+e}{1+e}$



15. A perfectly conducting fluid, of permittivity ε and permeability μ , flows with a uniform velocity \mathbf{v} in the presence of time dependent electric and magnetic fields \mathbf{E} and \mathbf{B} , respectively. If there is a finite current density in the fluid, then

(a) $\nabla \times (\mathbf{v} \times \mathbf{B}) = \frac{\partial \mathbf{B}}{\partial t}$ (b) $\nabla \times (\mathbf{v} \times \mathbf{B}) = -\frac{\partial \mathbf{B}}{\partial t}$
 (c) $\nabla \times (\mathbf{v} \times \mathbf{B}) = \sqrt{\varepsilon\mu} \frac{\partial \mathbf{E}}{\partial t}$ (d) $\nabla \times (\mathbf{v} \times \mathbf{B}) = -\sqrt{\varepsilon\mu} \frac{\partial \mathbf{E}}{\partial t}$

16. The energy levels of a non-degenerate quantum system are $\varepsilon_n = nE_0$, where E_0 is a constant and $n = 1, 2, 3, \dots$. At a temperature T , the free energy F can be expressed in terms of the average energy E by

(a) $E_0 + k_B T \ln \frac{E}{E_0}$ (b) $E_0 + 2k_B T \ln \frac{E}{E_0}$ (c) $E_0 - k_B T \ln \frac{E}{E_0}$ (d) $E_0 - 2k_B T \ln \frac{E}{E_0}$

17. The unnormalized wavefunction of a particle in one dimension in an infinite square well with walls at $x=0$ and $x=a$, is $\psi(x) = x(a-x)$. If $\psi(x)$ is expanded as a linear combination of the energy eigenfunctions, $\int_0^a |\psi(x)|^2 dx$ is proportional to the infinite series

[You may use: $\int_0^a t \sin t dt = -a \cos a + \sin a$ and $\int_0^a t^2 \sin t dt = -2 - (a^2 - 2) \cos a + 2a \sin a$]

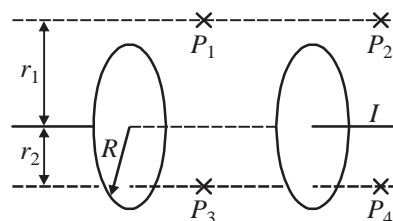
(a) $\sum_{n=1}^{\infty} (2n-1)^{-6}$ (b) $\sum_{n=1}^{\infty} (2n-1)^{-4}$ (c) $\sum_{n=1}^{\infty} (2n-1)^{-2}$ (d) $\sum_{n=1}^{\infty} (2n-1)^{-8}$

18. The Legendre polynomials $P_n(x)$, $n=0, 1, 2, \dots$, satisfying the orthogonality condition $\int_{-1}^1 P_n(x) P_m(x) dx = \frac{2}{2n+1} \delta_{nm}$ on the interval $[-1, +1]$, may be defined by the Rodrigues formula

$P_n(x) = \frac{1}{2^n n!} \frac{d^n}{dx^n} (x^2 - 1)^n$. The value of the definite integral $\int_{-1}^1 (4 + 2x - 3x^2 + 4x^3) P_3(x) dx$ is

(a) 3/5 (b) 11/15 (c) 23/32 (d) 16/35

19. The figure below shows an ideal capacitor consisting of two parallel circular plates of radius R . Points P_1 and P_2 are at a transverse distance $r_1 > R$ from the line joining the centres of the plates, while points P_3 and P_4 are at a transverse distance $r_2 < R$.

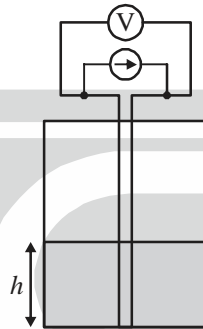


If $B(x)$ denotes the magnitude of the magnetic fields at these points, which of the following holds while the capacitor is charging ?

(a) $B(P_1) < B(P_2)$ and $B(P_3) < B(P_4)$ (b) $B(P_1) > B(P_2)$ and $B(P_3) > B(P_4)$
 (c) $B(P_1) = B(P_2)$ and $B(P_3) < B(P_4)$ (d) $B(P_1) = B(P_2)$ and $B(P_3) > B(P_4)$

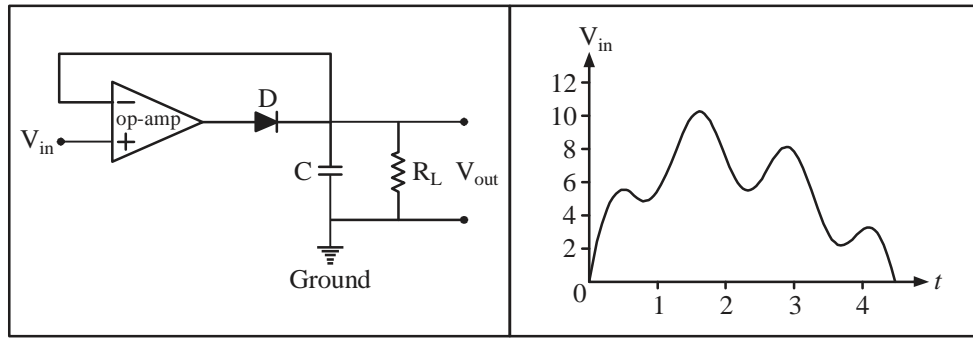
20. Balls of ten different colours labeled by $a = 1, 2, \dots, 10$ are to be distributed among different coloured boxes. A ball can only go in a box of the same colour, and each box can contain at most one ball. Let n_a and N_a denote, respectively, the numbers of balls and boxes of colour a . Assuming that $N_a \gg n_a \gg 1$, the total entropy (in units of the Boltzmann constant) can be best approximated by
- (a) $\sum_a (N_a \ln N_a + n_a \ln n_a - (N_a - n_a) \ln (N_a - n_a))$
 (b) $\sum_a (N_a \ln N_a - n_a \ln n_a - (N_a - n_a) \ln (N_a - n_a))$
 (c) $\sum_a (N_a \ln N_a - n_a \ln n_a + (N_a - n_a) \ln (N_a - n_a))$
 (d) $\sum_a (N_a \ln N_a + n_a \ln n_a + (N_a - n_a) \ln (N_a - n_a))$

21. To measure the height h of a column of liquid helium in a container, a constant current I is sent through an NbTi wire of length l , as shown in the figure. The normal state resistance of the NbTi wire is R .

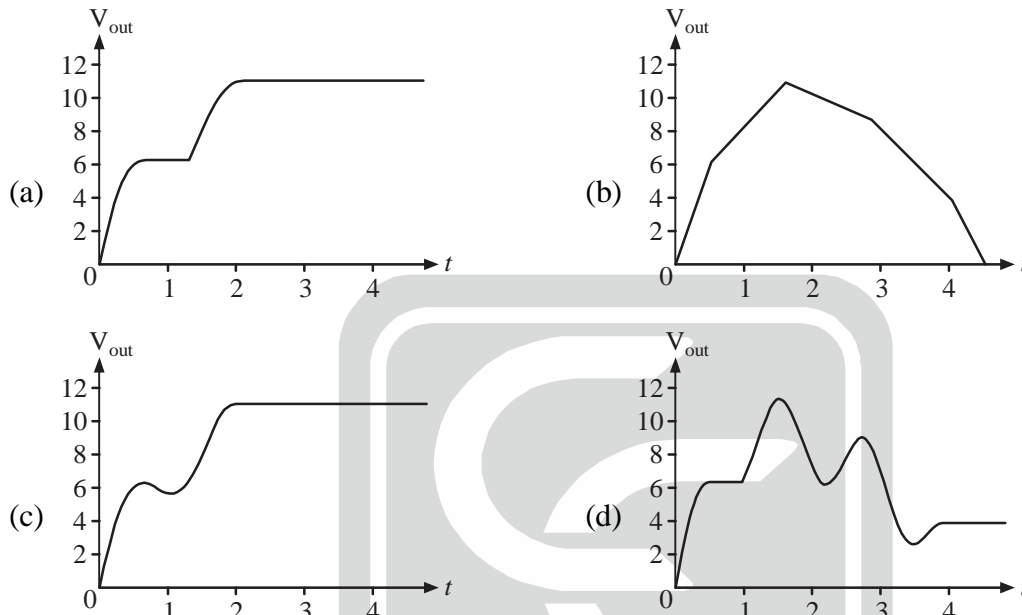


If the superconducting transition temperature of NbTi is ≈ 10 K, then the measured voltage $V(h)$ is best described by the expression.

- (a) $IR\left(\frac{1}{2} - \frac{2h}{l}\right)$ (b) $IR\left(1 - \frac{h}{l}\right)$ (c) $IR\left(\frac{1}{2} - \frac{h}{l}\right)$ (d) $IR\left(1 - \frac{2h}{l}\right)$
22. A particle of mass m in one dimension is in the ground state of a simple harmonic oscillator described by a Hamiltonian $H = \frac{p^2}{2m} + \frac{1}{2}m\omega^2 x^2$ in the standard notation. An impulsive force at time $t = 0$ suddenly imparts a momentum $p_0 = \sqrt{\hbar m \omega}$ to it. The probability that the particle remains in the original ground state is
- (a) e^{-2} (b) $e^{-3/2}$ (c) e^{-1} (d) $e^{-1/2}$
23. A particle in two dimensions is found to trace an orbit $r(\theta) = r_0 \theta^2$. If it is moving under the influence of a central potential $V(r) = c_1 r^{-a} + c_2 r^{-b}$, where r_0, c_1 and c_2 are constants of appropriate dimensions, the values of 'a' and 'b', respectively, are
- (a) 2 and 4 (b) 2 and 3 (c) 3 and 4 (d) 1 and 3
24. In the following circuit the input voltage V_{in} is such that $|V_{in}| < |V_{sat}|$, where V_{sat} is the saturation voltage of the op-amp. (Assume that the diode is an ideal one and $R_L C$ is much larger than the duration of the measurement)



For the input voltage as shown in the figure above, the output voltage V_{out} is best represented by



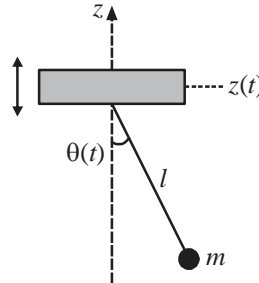
25. Lead is superconducting below 7 K and has a critical magnetic field 800×10^{-4} Tesla close to 0 K. At 2 K the critical current that flows through a long lead wire of radius 5 mm is closest to
 (a) 1760 A (b) 1670 A (c) 1950 A (d) 1840 A
26. If we use the Fourier transform $\phi(x, y) = \int e^{ikx} \phi_k(y) dk$ to solve the partial differential equation:

$$-\frac{\partial^2 \phi(x, y)}{\partial y^2} - \frac{1}{y^2} \frac{\partial^2 \phi(x, y)}{\partial x^2} + \frac{m^2}{y^2} \phi(x, y) = 0$$

in the half-plane $\{(x, y) : -\infty < x < \infty, 0 < y < \infty\}$ the Fourier modes $\phi_k(y)$ depend on y as y^α and y^β . The values of α and β are

- (a) $\frac{1}{2} + \sqrt{1 + 4(k^2 + m^2)}$ and $\frac{1}{2} - \sqrt{1 + 4(k^2 + m^2)}$
 (b) $1 + \sqrt{1 + 4(k^2 + m^2)}$ and $1 - \sqrt{1 + 4(k^2 + m^2)}$
 (c) $\frac{1}{2} + \frac{1}{2} \sqrt{1 + 4(k^2 + m^2)}$ and $\frac{1}{2} - \frac{1}{2} \sqrt{1 + 4(k^2 + m^2)}$
 (d) $1 + \frac{1}{2} \sqrt{1 + 4(k^2 + m^2)}$ and $1 - \frac{1}{2} \sqrt{1 + 4(k^2 + m^2)}$

27. The fulcrum of a simple pendulum (consisting of a particle of mass m attached to the support by a massless string of length l) oscillates vertically as $z(t) = a \sin \omega t$, where ω is a constant. The pendulum moves in a vertical plane and $\theta(t)$ denotes its angular position with respect to the z -axis.



If $l \frac{d^2\theta}{dt^2} + \sin \theta (g - f(t)) = 0$, (where g is the acceleration due to gravity) describes the equation of motion of the mass, then $f(t)$ is

- (a) $a\omega^2 \cos \omega t$ (b) $a\omega^2 \sin \omega t$ (c) $-a\omega^2 \cos \omega t$ (d) $-a\omega^2 \sin \omega t$
28. The Q -value of the α -decay of ^{232}Th to the ground state of ^{228}Ra is 4082 keV. The maximum possible kinetic energy of the α -particle is closest to
 (a) 4082 keV (b) 4050 keV (c) 4035 keV (d) 4012 keV
29. The nuclei of ^{137}Cs decay by the emission of β -particles with a half-life of 30.08 years. The activity (in units of disintegrations per second or Bq) of a 1 mg source of ^{137}Cs , prepared on January 1, 1980, as measured on January 1, 2021 is closest to
 (a) 1.79×10^{16} (b) 1.79×10^9 (c) 1.24×10^{16} (d) 1.24×10^9
30. In an elastic scattering process at an energy E , the phase shifts satisfy $\delta_0 \approx 30^\circ$, $\delta_1 \approx 10^\circ$, while the other phase shifts are zero. The polar angle at which the differential cross-section peaks is closest to
 (a) 20° (b) 10° (c) 0° (d) 30°