Chapter 4

Electrostatic Energy

Electrostatic Energy:

The work necessary to assemble a system of charges against coulomb forces is stored in the field as potential energy. This is known as electrostatic energy.

Electrostatic energy system of point charges:

Let at any time are place a point charge q_1 at \vec{r}_1 \overline{a} . To place the charge q_1 at the position \vec{r}_1 \rightarrow *,* we require no work because there is no interacting coulomb field.

$$
\therefore U_1 = 0
$$

But to bring the charge q_2 to the position \vec{r}_2 \rightarrow we require work done against coulomb repulsion due to q_1 .

This work

$$
U_2 = -\frac{1}{4\pi\varepsilon_0} \int_{\infty}^{r_2} \frac{q_1 q_2}{r^2} dr = \frac{1}{4\pi\varepsilon_0} \frac{q_1 q_2}{|r_1 - r_2|} = \frac{1}{4\pi\varepsilon_0} \frac{q_1 q_2}{r_{12}}
$$

Therefore for two point charges,

$$
u = \frac{q_1 q_2}{4\pi \varepsilon_0 r_{12}}
$$

To bring another charge, we need work done.

$$
U_3 = \frac{1}{4\pi\varepsilon_0} \left[\frac{q_1\mu_3}{r_{13}} + \frac{q_2q_3}{r_{23}} \right]
$$

Therefore, net work done

$$
U = U_1 + U_2 + U_3 = 0 + \frac{1}{4\pi\varepsilon_0} \left[\frac{q_1 q_2}{r_{12}} + \frac{q_1 q_3}{r_{13}} + \frac{q_2 q_3}{r_{23}} \right]
$$

For N point charged,

$$
U = \frac{1}{2} \sum_{i=1}^{N} \sum_{j=1}^{N} \frac{1}{4\pi \varepsilon_0} \frac{q_i q_j}{r_{ij}} \qquad \qquad q_i \qquad \qquad \sum_{r_{12}}^{r_{13}} \sum_{r_{12}}^{r_{23}}
$$

We know electrostatic potential

$$
\phi_i = \sum_{\substack{j=1\\i \neq j}}^N \frac{1}{4\pi\varepsilon_0} \frac{q_j}{r_{ij}}
$$

DEAVOI

 $\ddot{\cdot}$

$$
U=\frac{1}{2}\sum_{i=1}^N q_i\phi_i
$$

Potential energy of interaction of a charge 'q' with other charges

 $u = qV$, where *V* is potential at location of *q* due to other charges.

For continuous charge distribution :

For continous charge distribution u is given by

$$
U = \frac{1}{2} \int \rho(\vec{r}) \phi(\vec{r}) dr
$$
 [$\rho(\vec{r})$ is volume charge distribution]

Electrostatic Energy in terms of field :

Suppose we have a finite region of space V in a dielectric medium of permitivity ε and the volume charge density ρ .

Therefore, the electrostatic energy of the system is given by

$$
U = \frac{1}{2} \varepsilon_0 \int \rho(r) \phi(r) dV = \frac{1}{2} \int (\vec{\nabla} \cdot \vec{E}) \phi dV \qquad \text{[Gauss's law } \varepsilon_0 \vec{\nabla} \cdot \vec{E} = \rho \, \text{]}
$$

= $\frac{1}{2} \varepsilon_0 \int (\vec{\nabla} \cdot \phi \vec{E}) - \vec{\nabla} \phi \cdot \vec{E} \, dV$
= $\frac{1}{2} \varepsilon_0 \int (\vec{\nabla} \cdot \phi \vec{E}) dV - \frac{1}{2} \int (\vec{\nabla} \phi \cdot \vec{E}) dV$
= $\frac{1}{2} \varepsilon_0 \int (\phi \vec{E}) dS + \frac{1}{2} (\vec{E} \cdot \vec{E}) dV$

At large distance from the charge distribution the first term will be vanish.

$$
U = \frac{1}{2} \varepsilon_0 \int E^2 dV
$$

energy density, $u = \frac{1}{2} \varepsilon_0 E^2$ **ENDEAVOUR**

Therefore, the

Forces and Torques from electrostatic energy:

We know work done = force . displacement \vec{r} \vec{r}

$$
\Rightarrow \qquad dW = \vec{F} \cdot d\vec{r}
$$

Also for isolated system, the work done

$$
dW = -dU = -\vec{\nabla}U \cdot d\vec{r}
$$

$$
\vec{F} = -\vec{\nabla}U
$$
, $u = -\int \vec{F} \cdot d\vec{r}$

If the part of the system under consideration is constrained to rotate about an axis and it undergoes a small rotation do under the action of a torque $\vec{\tau}$ then,

$$
dW = \vec{\tau} \cdot d\vec{\theta} = -dU = -\nabla_{\theta}U \cdot d\theta
$$

$$
\vec{\tau} = -\nabla_{\theta}U \quad \text{and } u = \int \vec{F} \cdot d\vec{r}
$$

 \mathcal{L}

The electrostatic energy associated with a uniform spherical charge distribution of total charge *q* **and radius** *R***.**

The electric field of this distribution is

$$
E = E_1 = \frac{qr}{4\pi\varepsilon_0 R^3} \text{ for } r \le R
$$

$$
E = E_2 = \frac{q}{4\pi\varepsilon_0 R^2} \text{ for } r \ge R
$$

Since, the field is radially symmetric, the volume element may be expressed as $dv = 4\pi r^2 dr$, so that the energy is

$$
U = \frac{\varepsilon_0}{2} \int E^2 dv = \frac{\varepsilon_0}{2} \int_0^R E_1^2 4\pi r^2 dr + \frac{\varepsilon_0}{2} \int_R^{\infty} E_2^2 4\pi r^2 dr
$$

$$
= \frac{1}{2} \int_0^R \frac{q^2 r^2}{16\pi^2 \varepsilon_0 R^6} 4\pi r^2 dr + \frac{1}{2} \int_R^{\infty} \frac{q^2}{16\pi^2 \varepsilon_0 r^4} 4\pi r^2 dr
$$

$$
= \frac{q^2}{40\pi \varepsilon_0 R} + \frac{q^2}{8\pi \varepsilon_0 R}, \quad U = \frac{3q^2}{20\pi \varepsilon_0 R}
$$

The electrostatic energy of a uniformly charged spherical shell of total charge *q* and radius *a*.

The electric field of the shell is

$$
E = 0, \qquad \text{for } r \le R
$$

$$
E = \frac{q}{4\pi e_0 r^2} \qquad \text{for } r \ge R
$$

The energy is therefore,

$$
U = \frac{\varepsilon_0}{2} \int E^2 dv = \frac{1}{2} \int_{R}^{\infty} \frac{q^2}{16\pi^2 \varepsilon_0 r^4} 4\pi r^2 dr \quad \text{DEAVOLR}
$$

$$
U=\frac{q^2}{8\pi\varepsilon_0 R}
$$

Example: Consider a spherical charge distribution with a volume density

$$
\rho(r) = \begin{cases} \rho_0 \cdot \frac{a}{r} & \text{for } 0 < r < a \\ 0 & \text{for } r > a \end{cases}
$$

Show that the total electrostatic energy of the system is 2 $6\pi \in _{0}$ *Q* $\sum_{\pi \in \Omega}$ where Q is the total charge in the

system.

Soln. Total charge in the system is

$$
Q = \int_{0}^{a} \rho(r) \cdot 4\pi r^2 dr = 4\pi \rho_0 a \int_{0}^{a} r dr = 2\pi \rho_0 a^3
$$

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$$
E = \frac{Q_{enclosed}}{4\pi\varepsilon_0 r^2}
$$
 or $Q_{enclosed} = \int \rho d\tau = 4\pi\rho_0 a \int_0^a r dr = 2\pi\rho_0 a^3$

 \sum

For an internal point,

$$
\int \vec{E}_i \cdot d\vec{S} = \frac{Q_r}{\epsilon_0} = \frac{1}{\epsilon_0} \int_0^r \rho(r) \cdot 4\pi r^2 dr = \frac{2\pi \rho_0 a r^2}{\epsilon_0}
$$

$$
\therefore E_i = \frac{Q_{enclosed}}{4\pi\varepsilon_0 r^2} \text{ or } E_i = \frac{Q}{4\pi\varepsilon_0 a^2}
$$

For external point, $E_{out} = \frac{2 \times 10^2}{4 \pi \epsilon_0 r^2} = \frac{2 \times 10^4}{4 \pi \epsilon_0 r^2}$ $\frac{Q_{enclosed}}{A - \alpha}$ $E_{\textit{\tiny out}} = \frac{Q_{\textit{enclosed}}}{\frac{1}{2} - \frac{Q}{A}} = \frac{Q}{A}$ $\pi \varepsilon_0 r^2$ $4\pi \varepsilon_0 r^2$ $=\frac{\sum_{enclosed}}{1}$ = -

Now, the total energy in the system is

$$
U = \int_{0}^{a} \frac{1}{2} \epsilon_0 E_i^2 \cdot 4\pi r^2 dr + \int_{a}^{\infty} \frac{1}{2} \epsilon_0 E_{out}^2 \cdot 4\pi r^2 dr
$$

= $2\pi \epsilon_0 \int_{0}^{a} \left(\frac{Q}{4\pi \epsilon_0 a^2} \right)^2 r^2 dr + 2\pi \epsilon_0 \int_{a}^{\infty} \left(\frac{Q}{4\pi \epsilon_0} \right)^2 \cdot \frac{dr}{r^2} = \frac{Q^2}{6\pi \epsilon_0 a}$

ADDITIONAL SOLVED PROBLEMS

- 1. A particle of mass 40 *mg* and carrying a charge 5×10^{-9} C is moving directly towards a fixed positive point charge of magnitude 10^{-8} C. When it is at a distance of 10 cm from the fixed positive point charge it has a velocity of 50 cm sec–1. At what distance from the fixed point charge will the particle come momentarily to rest? Is the acceleration constant during motion ?
- **Soln.** Let *x* be the distance of q from *Q* when it comes to rest as shown in figure:

$$
Q = 10^{-8} \text{ C}
$$

\n
$$
v = 0
$$

\n
$$
m = 40 \times 10^{-6} \text{ kg}
$$

\n
$$
u = 0.5 \text{ m/sec}
$$

\n
$$
v = 0
$$

Use conservation of energy $(K.E. + P.E.)$ _{initial} $=(K.E. + P.E.)$ _{final}

$$
\frac{1}{2} \times 40 \times 10^{-6} \times (0.5)^2 + \frac{10^{-8} \times 5 \times 9 \times 10^9}{\left(\frac{1}{10}\right)} = 0 + \frac{10^{-8} \times 5 \times 10^{-9} \times 9 \times 10^9}{x}
$$

Solving we get, $r = 4.737 \times 10^{-2}$ m.

The force obeys inverse square law i.e., it increases as distance decreases. Hence, acceleration also increases.

- 2. Four charges $+q$, $+q$, $-q$ and $-q$ are placed respectively at the corners *A*, *BC* and *D* of a square of side *a*, arranged in the given order. If *E* and *F* are mid points of side *BC* and *CD* respectively, what will be workdone in carrying a charge '*e*' from *O* to *E* and from *O* to *F*.
- **Soln.** The charge configuration is shown in figure :

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Workdone in carrying a charge = change in potential energy = $q(V_{final} - V_{initial})$

$$
AE^{2} = a^{2} + \left(\frac{a}{2}\right)^{2} = \frac{5a^{2}}{4} \text{ or } AE = \frac{\sqrt{5} \cdot a}{2}
$$

\nPotential at *E* i.e., $V_{E} = \frac{1}{4\pi \epsilon_{0}} \left[\frac{q}{\sqrt{5} \cdot \frac{a}{2}} - \frac{q}{\sqrt{5} \cdot \frac{a}{2}} + \frac{q}{a/2} - \frac{q}{a/2} \right] = 0$
\nPotential at *F* i.e., V_{F} is given by
\n
$$
V_{F} = \frac{1}{4\pi \epsilon_{0}} \left[\frac{q}{\sqrt{5} \cdot \frac{a}{2}} + \frac{q}{\sqrt{5} \cdot \frac{a}{2}} - \frac{q}{a/2} - \frac{q}{a/2} \right]
$$
\n
$$
= \frac{1}{4\pi \epsilon_{0}} \left[\frac{4q}{\sqrt{5} \cdot a} - \frac{4q}{a} \right]
$$
\n
$$
= \frac{1}{4\pi \epsilon_{0}} \cdot \frac{4q}{a} \left[\frac{1}{\sqrt{5}} - 1 \right]
$$
\nNow, $W_{OE} = e(V_{E} - V_{O}) = e(0 - 0) = 0$ and $W_{OF} = e(V_{F} - V_{O})$
\n
$$
= e \cdot \left[\frac{1}{4\pi \epsilon_{0}} \cdot \frac{4q}{a} \left\{ \frac{1}{\sqrt{5}} - 1 \right\} - 0 \right] = \frac{1}{4\pi \epsilon_{0}} \cdot \frac{4eq}{a} \cdot \left[\frac{1}{\sqrt{5}} - 1 \right].
$$

3. Two fixed, equal positive charges, each of magnitude 5×10^{-5} coul, are located at points *A* and *B* separated at a distance of 6 m. An equal and opposite charge moves towards them along the line *COD*, the perpendicular bisector of the line *AB*. The moving charge, when reaches the point *C* at a distance of 4 m from *O*, has a kinetic energy of 4 joules. Calculate the distance of the farthest point *D* at which the negative charge will reach before returning towards *C*.

Soln. Applying conservation of energy,
$$
(K.E.+P.E.)_{\text{arc}} = (K.E.+P.E.)_{\
$$

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- 4. Three charges, q , $2q$ and $8q$ are to be placed on 9 cm long straight line. Find the positions where the charges should be placed such that the potential energy of this system is minimum.
- **Soln.** For minimum interaction energy, the charges of greater magnitudes should be placed at the extreme ends. The situation is shown in figure. Let the distance of charge 2*q* from charge *q* be *x* cm.

The potential energy of the system is given by

$$
U = \frac{1}{4\pi\varepsilon_0} \left[\frac{2q \times q}{x} + \frac{q \times 8q}{(9-x)} + \frac{2q \times 8q}{9} \right]
$$

For minimum potential energy, $\partial U / \partial x = 0$

$$
\therefore \qquad \frac{\partial U}{\partial x} = \frac{1}{4\pi\epsilon_0} \Rightarrow \frac{\partial}{\partial x} \left[\frac{2q^2}{x \times 10^{-2}} + \frac{8q^2}{(9-x)10^{-2}} + \frac{16q^2}{9 \times 10^{-2}} \right] = 0
$$

Solving we get $x = 3$ cm.

Thus the charge *q* should be placed at a distance of 3 cm from the charge 2*q*.

- 5. A circular ring of radius *R* with uniform positive charge density λ per unit length is located in the YZ-plane with its centre at the origin *O*. A particle of mass *m* and positive charge *q* is projected from the point $P(R\sqrt{3}, 0, 0)$ on the positive *X*-axis directly towards *O*, with initial velocity *v*. Find the smallest (non-zero) value of the speed *v* such that the particle does not return to *P*.
- **Soln.** The situation is shown in figure :

Particle will not return to P if it manages to at least reach center of ring. Potential at axial point of ring (charge

$$
= Q, \text{ radius} = R) \text{ is } V_{\text{ring}} = \frac{Q}{4\pi\varepsilon_0\sqrt{R^2 + x^2}}
$$

Potential energy of charge and ring = qV_{ring}

Appling conservation of energy

$$
(K.E. + P.E.)_{x=\sqrt{3} R} = (K.E. + P.E.)_{x=0}
$$

$$
\frac{1}{2}mv^2 + \frac{qQ}{4\pi\varepsilon_0\sqrt{R^2 + (\sqrt{3}R)^2}} = 0 + \frac{qQ}{4\pi\varepsilon_0\sqrt{R^2 + 0^2}} \Rightarrow v = \sqrt{\frac{qQ}{4\pi\varepsilon_0 R m}}, \text{ since } Q = \lambda 2\pi R
$$

$$
\therefore \qquad v = \frac{\lambda q}{2\varepsilon_0 m}
$$

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- 6. A conducting sphere S_1 , of radius *r* is attached to an insulating handle. Another conducting sphere S_2 of radius *R* is mounted as an insulating stand. S_2 is initially uncharged. S_1 is given a charge Q , brought into contact with S_2 and removed. S_1 is recharged such that the charge on it is again Q and is again brought into contact with
	- 2 *S* and removed. This procedure is repeated *n* times.
	- (a) Find the electrostatic energy of S_2 after *n* such contacts with S_1 .
	- (b) What is the limiting value of this energy as $n \to \infty$.
- **Soln.** (a) Let Q_2 be the charge on sphere S_2 after S_1 is removed first Q_1 be charge on S_1 .

$$
\therefore \qquad Q_1 + Q_2 = Q
$$

Since the potential at the surface of each sphere is the same, we have

$$
\frac{Q_2}{4\pi\varepsilon_0 R} = \frac{Q_1}{4\pi\varepsilon_0 r} \Rightarrow \frac{Q_2}{R} = \frac{Q - Q_2}{r} \text{ i.e. } Q_2 = \frac{1}{1 + \frac{r}{R}}
$$
 times, the total charge $Q_2 = \frac{Q}{1 + \left(\frac{r}{R}\right)}$

Now S_1 is recharged such that the charge on it is again *Q*. Charge on S_2 is $\left(1+\frac{r}{R}\right)$ *R Q* $^{+}$. They are brought in

contact.

Total charge on
$$
S_1
$$
 and $S_2 = \left[Q + \left\{ \frac{Q}{\left(1 + \frac{r}{R} \right)} \right\} \right]$

Therefore, now total charge on $Q_{\scriptscriptstyle 2}$ will be

$$
= \frac{1}{1+\frac{r}{R}}
$$
time total charge
$$
= \frac{1}{\left(1+\frac{r}{R}\right)} \cdot \left[Q + \frac{Q}{1+\frac{r}{R}}\right] = \frac{Q}{1+\frac{r}{R}} + \frac{Q}{\left(1+\frac{r}{R}\right)^{2}}
$$

Similarly, after *n* repeatation, total charge on S_2 becomes,

$$
= \frac{Q}{\left\{1+\left(\frac{r}{R}\right)\right\}} + \frac{Q}{\left\{1+\left(\frac{r}{R}\right)\right\}^2} + \dots + \frac{Q}{\left\{1+\left(\frac{r}{R}\right)\right\}^n}
$$

$$
= \frac{Q}{1+\left(\frac{r}{R}\right)} \left[\frac{1-1/(1+\frac{r}{R})^n}{1-1/(1+\frac{r}{R})} \right]
$$

Electrostatic energy of S_2

$$
= \frac{Q_{\text{final}}^2}{8\pi\epsilon_0 R} = \frac{Q^2}{8\pi\epsilon_0 R \left(1 + \frac{r}{R}\right)^2} \times \left[\frac{1 - 1/(1 + \frac{r}{R})^n}{1 - 1/(1 + \frac{r}{R})}\right]^2
$$

(b) Limiting energy at $n \to \infty$

$$
E = \frac{Q^2}{8\pi\epsilon_0 R \left(1 + \frac{r}{R}\right)^2} \left[\frac{1 - 0}{1 - 1/(1 + \frac{r}{R})}\right]^2
$$

$$
Q^2 = (R)^2 \cdot Q^2 R
$$

$$
=\frac{Q^2}{8\pi\epsilon_0 R}\left(\frac{R}{r}\right)^2=\frac{Q^2 R}{8\pi\epsilon_0 r^2}.
$$

- 7. A non-conducting disc of radius *a* and uniform surface charge density σ is placed on the ground, with its axis vertical. A particle of mass *m* and positive charge *q* is dropped, along the axis of the disc, from a height *H* with zero initial velocity. The particle has $q/m = 4\pi \varepsilon_0 g/\sigma$.
	- (a) Find the value of H if the particle just reaches the disc.
	- (b) Sketch the potential energy of the particle as a function of its height and find its equilibrium position.
- **Soln.** See below figure.

We know that, $U = qV_{disc}$

$$
=\frac{\sigma q}{2\epsilon_0} \left[(a^2 + x^2)^{1/2} - x \right] \qquad \qquad \dots (1)
$$

(a) Applying the principle of conservation of energy at points P and O, we have $(K.E. + P.E.)_{at P} = (K.E. + P.E.)_{at O}$

$$
0+m g H+\frac{\sigma q}{2\epsilon_0}\left[\sqrt{\left(a^2+H^2\right)-H}\right]=0+\frac{\sigma q}{2\epsilon_0}\left[\sqrt{\left(a^2+0^2\right)-0}\right]
$$

 ε_0 g m/ σ) Given that $q = (4 \varepsilon_0 g m/\sigma)$

$$
\therefore \quad mg \, H + \frac{\sigma(4\,\varepsilon_0 \, g \, m/\sigma)}{2\,\varepsilon_0} \bigg[\sqrt{(a^2 + H^2)} - H \bigg] = \frac{\sigma(4\pi\,\varepsilon_0 \, g \, m/\sigma)}{2\,\varepsilon_0} a
$$

or
$$
mg \, H + 2mg \bigg[\sqrt{(a^2 + H^2)} - H \bigg] = 2m \, g \, a
$$

$$
H + 2\sqrt{(a^2 + H^2)} - 2H = 2a
$$

$$
2\sqrt{(a^2 + H^2)} = (2a + H)
$$

or
$$
4(a^2 + H^2) = (2a + H)^2
$$

Solving we get $H = (4a/3)$

Solving we get $H = (4a/3)$

(b) Total potential energy at a height *x* above centre

$$
P.E. = \frac{\sigma q}{2 \varepsilon_0} \left[\sqrt{(a^2 + x^2)} - x \right] + mg \ x
$$

 $2 m g a$

3*^m ga*

P.E.

Substituting the value of *q*, we get

$$
P.E. = 2mg\left[\sqrt{(a^2 + x^2)} - x\right] + mg x
$$

$$
U = mg\left[2\sqrt{(a^2 + x^2)} - x\right]
$$

For equilibrium position, $\frac{dU}{dt} = 0$ *dx* $=$

$$
\therefore \qquad \frac{dU}{dx} = 0 = mg \left[\frac{2.2x}{2\sqrt{(a^2 + x^2)}} - 1 \right]
$$

or
$$
\frac{2x}{\sqrt{(a^2 + x^2)}} = 1
$$
 or $4x^2 = (a^2 + x^2)$

$$
\therefore \qquad x = \frac{a}{\sqrt{3}}
$$

The variation of potential energy with x is shown in figure.

